

HΦの概要

三澤 貴宏

東京大学物性研究所計算物質科学研究センター



<http://www.pasums.issp.u-tokyo.ac.jp/hphi/>

1. Introduction

Basics of models for SCES

Why do we want to solve models for SCES ?

How difficult to solve models for SCES ?

2. Algorithms used in $H\Phi$

- exact diagonalization (Lanczos, LOBCG)
- finite-temperature calculations

3. How to use $H\Phi$

Introduction



What is *models* for SCES ?

e.g. Hubbard model

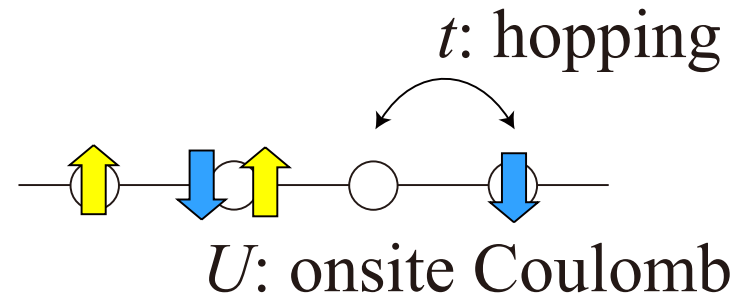
$$\hat{\mathcal{H}} = \hat{\mathcal{H}}_t + \hat{\mathcal{H}}_U$$

Electrons as waves

$$\hat{H}_t = -t \sum_{\langle i,j \rangle, \sigma} (\hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + \hat{c}_{j\sigma}^\dagger \hat{c}_{i\sigma})$$

Electrons as particles

$$\hat{H}_U = \sum_i \hat{n}_{i\uparrow} \hat{n}_{i\downarrow}, \quad \hat{n}_{i\sigma} = \hat{c}_{i\sigma}^\dagger \hat{c}_{i\sigma}$$



Relations between 2nd-quantized operators (these are all !)

$$\{\hat{c}_{i\sigma}^\dagger, \hat{c}_{j\sigma'}\} = \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma'} + \hat{c}_{j\sigma'} \hat{c}_{i\sigma}^\dagger = \delta_{i,j} \delta_{\sigma,\sigma'}$$

$$\{\hat{c}_{i\sigma}^\dagger, \hat{c}_{j\sigma'}^\dagger\} = 0 \rightarrow \hat{c}_{i\sigma}^\dagger \hat{c}_{i\sigma}^\dagger = 0 \quad \text{Pauli's principle}$$

$$\{\hat{c}_{i\sigma}, \hat{c}_{j\sigma'}\} = 0 \rightarrow \hat{c}_{i\sigma} \hat{c}_{i\sigma} = 0$$

Why do we want to solve ?

Rules (=models) are known, but results are unknown

Japanese chess (将棋)

We know how each piece moves,



but we don't know which one wins.



Solids state physics

We know how each electron moves,

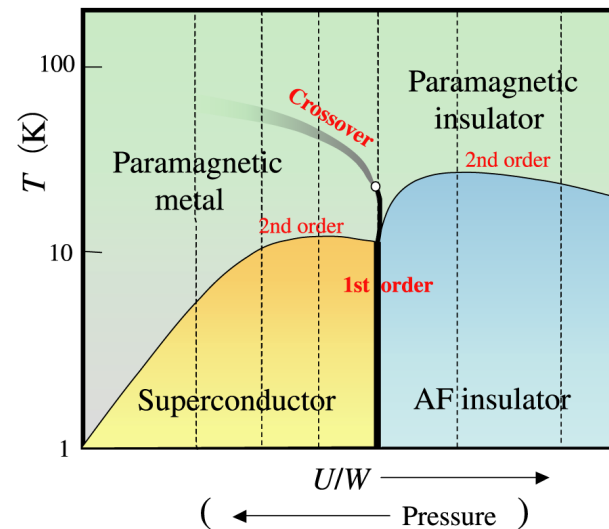
$$\hat{\mathcal{H}} = \hat{\mathcal{H}}_t + \hat{\mathcal{H}}_U$$

$$\{\hat{c}_{i\sigma}^\dagger, \hat{c}_{j\sigma'}\} = \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma'} + \hat{c}_{j\sigma'} \hat{c}_{i\sigma}^\dagger = \delta_{i,j} \delta_{\sigma,\sigma'}$$

$$\{\hat{c}_{i\sigma}^\dagger, \hat{c}_{j\sigma'}^\dagger\} = 0 \rightarrow \hat{c}_{i\sigma}^\dagger \hat{c}_{i\sigma}^\dagger = 0$$

$$\{\hat{c}_{i\sigma}, \hat{c}_{j\sigma'}\} = 0 \rightarrow \hat{c}_{i\sigma} \hat{c}_{i\sigma} = 0$$

but we don't know what happens.



K. Kanoda, JPSJ 75, 051007 (2006)

Wave function = eigenvectors of Hamiltonian

Matrix representation of Hamiltonian (ex. 2 site Hubbard model)

Real-space configuration $|\uparrow, \downarrow\rangle = c_{1\uparrow}^\dagger c_{2\uparrow}^\dagger |0\rangle$

After some *tedious* calculations, $\langle \uparrow, \downarrow | \hat{H}_t | \uparrow, \downarrow, 0 \rangle = \langle \uparrow, \downarrow | (t \sum_{\sigma} c_{1\sigma}^\dagger c_{2\sigma} + c_{2\sigma}^\dagger c_{1\sigma}) | \uparrow, \downarrow, 0 \rangle = -t$

$$\mathcal{H} = \begin{array}{c} \langle \uparrow, \downarrow | \\ \langle \downarrow, \uparrow | \\ \langle \uparrow\downarrow, 0 | \\ \langle 0, \uparrow\downarrow | \end{array} \begin{array}{cccc} | \uparrow, \downarrow \rangle & | \downarrow, \uparrow \rangle & | \uparrow\downarrow, 0 \rangle & | 0, \uparrow\downarrow \rangle \\ \left(\begin{array}{cccc} 0 & 0 & -t & -t \\ 0 & 0 & t & t \\ -t & t & U & 0 \\ -t & t & 0 & U \end{array} \right) \end{array}$$

$$|\phi\rangle = a_0 |\uparrow, \downarrow\rangle + a_1 |\downarrow, \uparrow\rangle + a_2 |\uparrow\downarrow, 0\rangle + a_3 |0, \uparrow\downarrow\rangle$$

Diagonalization \rightarrow eigenvalues, eigenvectors
 \rightarrow Problem is completely solved ($\mathbf{H}\Phi$)

How difficult to solve models ?

“nightmare” for theorists



- Exact diagonalization $\sim 4^{N_s} \times 4^{N_s}$

→ Only for small sizes [HΦ]

- Unbiased Quantum Monte Carlo method
[DSQSS]

→ Negative sign problem

- Many methods have been developed
DMFT (∞ dimensions \rightarrow 2,3 dimensions) [DCore]
DMRG (1 dimension \rightarrow 2 dimensions)
tensor-network (1,2 dimensions) [TeNES]
VMC (wide applicable range, but ..) [mVMC]

No all-round method exists

For small systems, using exact diagonalizations, we can obtain *exact* results.

Introduction
to $H\Phi$



Developers of H Φ (2015-)

M. Kawamura

T. Misawa

K. Yoshimi



Y. Yamaji



S. Todo



N. Kawashima

K. Ido



Development of H Φ is supported by
“*Project for advancement of software
usability in materials science*” by ISSP

Basics of HΦ



What can we do by HΦ?

For Hubbard model, **spin- S** Heisenberg model, Kondo-lattice model ...

- Full diagonalization [Scalapack & GPGPU]
- Ground state calculations by Lanczos method
- Finite-temperature calculations by thermal pure quantum (TPQ) states
- Low-energy states calculation by LOBCG method
- Dynamical properties [$S(q, \omega)$], real-time evo.

maximum system sizes@ ISSP system B (sekirei)

- spin 1/2: ~ 40 sites (S_z conserved)
- Hubbard model: ~ 20sites (# of particles & S_z conserved)

Meaning of name & logo



- Multiplying H to Φ ($H\Phi$)
- The cat means wave function in two ways
cat is a symbol of superposition.. (Schrödinger's cat)

Meaning of name & logo



- Multiplying H to Φ ($H\Phi$)
- The cat means wave function in two ways
cat is a symbol of superposition.. (Schrödinger's cat)

wake



sleep



$\Phi =$

+

Algorithms used in $H\Phi$!

The image shows a stylized black silhouette of the GitHub Octocat logo. The logo is a cat-like creature with a large head and a small body. The letters 'H' and 'Φ' are placed on the cat's face, serving as eyes. The entire logo is enclosed within a thick black circular outline.

$H\Phi$

Full diagonalization

Matrix representation of Hamiltonian

(real space basis) \rightarrow Full diagonalization for the matrix

$$H_{ij} = \langle i | \hat{H} | j \rangle \quad |i\rangle \text{ real-space basis}$$

ex. 2-site Hubbard $\mathcal{H} = \begin{pmatrix} 0 & 0 & -t & -t \\ 0 & 0 & t & t \\ -t & t & U & 0 \\ -t & t & 0 & U \end{pmatrix}$

dim. of matrix =

of real-space bases = exponentially large

ex. spin 1/2 system: $S_z=0$

$$N_s \binom{N_s}{2}$$

- $N_s=16$: dim.=12800,

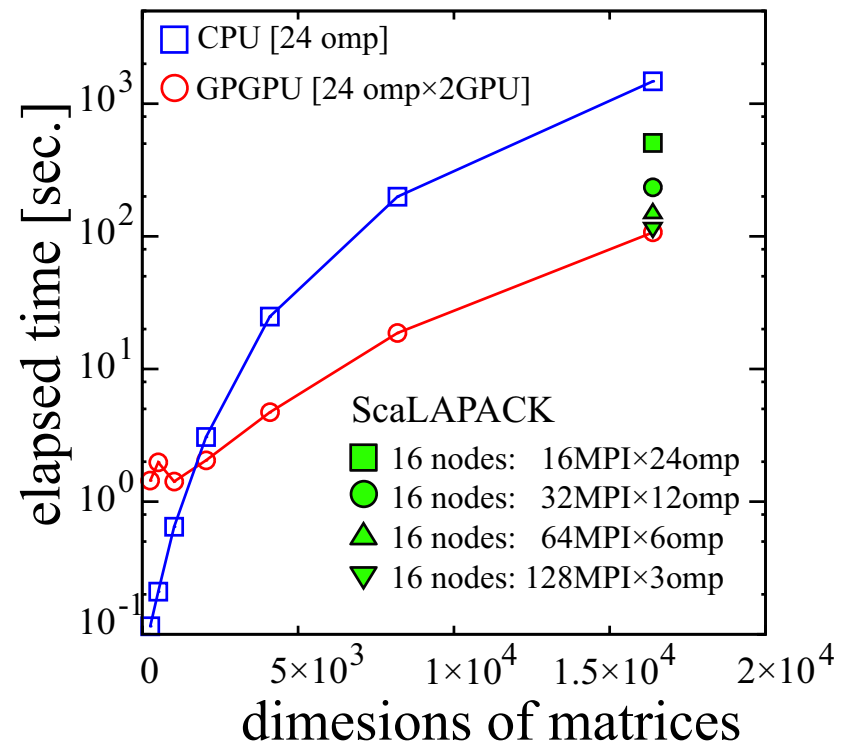
- $N_s=32$: dim. $\sim 6 \times 10^8$

required memory ~ 1 GB

required memory ~ 3 E B !

Full diagonalization

In $H\Phi$, full diagonalization can be done using LAPACK, SCALAPACK, MAGMA(GPGPU)



Input matrix can be specified by **MATRIX MARKET** format

```
%%%% matrix coordinate complex hermitian
```

```
4 4 6
```

```
1 1 4.000000 0.000000
2 1 1.000000 0.000000
3 1 -1.000000 0.000000
4 2 1.000000 0.000000
4 3 -1.000000 0.000000
4 4 4.000000 0.000000
```



A visual repository of test data for use in comparative studies of algorithms for numerical linear algebra, featuring nearly 500 sparse matrices from a variety of applications, as well as matrix generation tools and services.

Lanczos method

By multiplying the Hamiltonian to initial vector, we can obtain the ground state (power method)

$$\mathcal{H}^n \mathbf{x}_0 = E_0^n \left[a_0 \mathbf{e}_0 + \sum_{i \neq 0} \left(\frac{E_i}{E_0} \right)^n a_i \mathbf{e}_i \right]$$



A few (at least two) **vectors** are necessary →

We can treat larger system size than full diagonalization

ex. spin 1/2 system: $S_z=0$

- $N_s=16$: dim. = 12800, required memory (vectors) ~0.1 MB
- $N_s=32$: dim. $\sim 6 \times 10^8$, required memory (vectors) **~5 GB !**
- $N_s=36$: dim. $\sim 9 \times 10^9$, required memory (vectors) **~72 GB !**

LOBCG method

Lanczos method + conjugate gradient (CG) method

[Conventional exact diagonalization (ex. TITPACK)]

Lanczos[eigenvalues, 2 vectors] → Lanczos[eigenvectors, 3 vectors]

→ Inverted iteration [5 vectors]

[merit] Only two vectors are necessary

[demerit] Redundant, *Failure* of orthogonality, degeneracy, etc...

Locally Optimal (Preconditioning) Block Conjugate [LO(P)BCG]

Eigenvectors $\{|\Phi_0\rangle, |\Phi_1\rangle, \dots, |\Phi_{M-1}\rangle\}$

residual vectors $\{|r_i\rangle = H|\Phi_i\rangle - \epsilon_i|\Phi_i\rangle\}$

conjugate gradient vectors $\{|p_0\rangle, |p_1\rangle, \dots, |p_{M-1}\rangle\}$

A. V. Knyazev, SIAM J. Sci. Compute. 23, 517 (2001)

B. 山田進 他, 日本計算工学会論文集, 20060027 (2006)

Diagonalize & Optimize
with 3M vectors

[merit]

- Obtaining many low-energy states at once

- Degeneracy & orthogonality can be well reproduced

[demerit] Required memory is slightly larger (~3M-6M vectors)

Finite-temperature calc.!



Finite-temperature calculations

- Is ensemble average necessary for finite-temperature calculations ? **No !**

$$\langle E \rangle = \frac{\sum_n E_n e^{-\beta E_n}}{\sum_n e^{-\beta E_n}}$$

- Is it possible to replace ensemble average by the average values of wave function ? **(Partially) Yes !**

$$\langle E(\beta) \rangle = \frac{\langle \Phi(\beta) | H | \Phi(\beta) \rangle}{\langle \Phi(\beta) | \Phi(\beta) \rangle} ? \quad |\Phi(\beta)\rangle = \sum_i a_i(\beta) |i\rangle$$

Exact: Double Hilbert space (二重ヒルベルト空間, 熱場ダイナミクス)

高橋康, 物性研究 20, 97 (1973); 鈴木増雄, 統計力学(岩波書店)

Yasushi TAKAHASHI and Hiroomi UMEZAWA, *Int. J. Mod. Phys. B* 10, 1755 (1996).

Reprinted with permission from *Collective Phenomena* 1975, Vol. 2 pp. 55–80

Stochastic: Quantum Transfer Monte Carlo (Imada-Takahashi 1986)

Finite-temperature Lanczos (Jaklic-Prelovsek 1994)

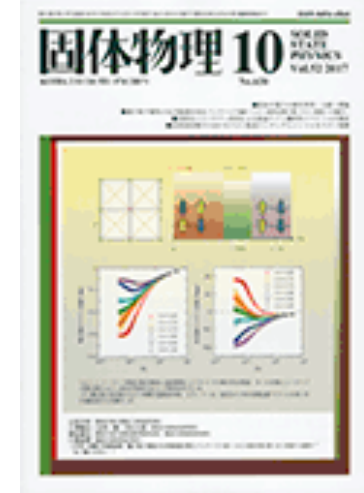
Hams-De Raedt method (2000)

Thermal Pure Quantum [TPQ] state (Sugiura-Shimizu 2012,2013)

Essence of TPQ

cf. 固体物理, vol. 52, No. 10

「量子格子模型の汎用数値対角化パッケージHΦ
—スピン液体近傍の熱・スピン励起への適用—」



1. Random vector (high-temperature limit) *equally* includes all eigenvectors

$$|\Phi_{\text{rand}}\rangle = \sum_n a_n |n\rangle, |a_n|^2 \sim 1/N_H$$

$$\langle \hat{H} \rangle = \frac{\langle \Phi_{\text{rand}} | \hat{H} | \Phi_{\text{rand}} \rangle}{\langle \Phi_{\text{rand}} | \Phi_{\text{rand}} \rangle} = \frac{\sum_n |a_n|^2 E_n}{\sum_n |a_n|^2} \sim \frac{1}{N_H} \sum_n E_n$$

2. Commutative quantities can be calculated by *single wave function*

$$|\Phi(\beta)\rangle = e^{-\beta \hat{H}/2} |\Phi_{\text{rand}}\rangle$$

$$\frac{\langle \Phi(\beta) | \hat{H} | \Phi(\beta) \rangle}{\langle \Phi(\beta) | \Phi(\beta) \rangle} = \frac{\sum_n |a_n|^2 E_n e^{-\beta E_n}}{\sum_n |a_n|^2 e^{-\beta E_n}} \sim \frac{\sum_n E_n e^{-\beta E_n}}{\sum_n e^{-\beta E_n}} = E(\beta)$$

3. Non-commutative quantities can be also calculated by *single wave function*

Proofs: Hams and De Raedt PRE **2000**; Sugiura and Shimizu PRL **2012,2013**
Thermal Pure Quantum state (熱的純粹量子状態) by Sugiura and Shimizu

Sugiura-Shimizu method [TPQ method]

Procedure

Simple implementation

$|\psi_0\rangle$: random vector

$$|\psi_k\rangle \equiv \frac{(l - \hat{H}/N_s)|\psi_{k-1}\rangle}{|(l - \hat{H}/N_s)|\psi_{k-1}\rangle|}$$

$$u_k \sim \langle \psi_k | \hat{H} | \psi_k \rangle / N_s$$

$$\beta_k \sim \frac{2k/N_s}{(l - u_k)}, \quad \langle \hat{A} \rangle_{\beta_k} \sim \langle \psi_k | \hat{A} | \psi_k \rangle$$

S. Sugiura and A. Shimizu,
PRL 2012

**l : constant larger the
maximum eigenvalues**

**All the finite temperature properties can be calculated
by using *one* thermal pure quantum [TPQ] state.**

Errors of estimation are bounded by S

→ Powerful tools for studying frustrated systems !

Error of estimation $\sim \exp[-S(T^*)/2]$ ($T/2 < T^* < T$)

Drastic reduction of numerical cost

Heisenberg model, 32 sites, $S_z=0$

Full diagonalization:

Dimension of Hamiltonian $\sim 10^8 \times 10^8$

Memory $\sim 3\text{E Byte}$ \rightarrow Almost impossible.

TPQ method:

Only two vectors are required:

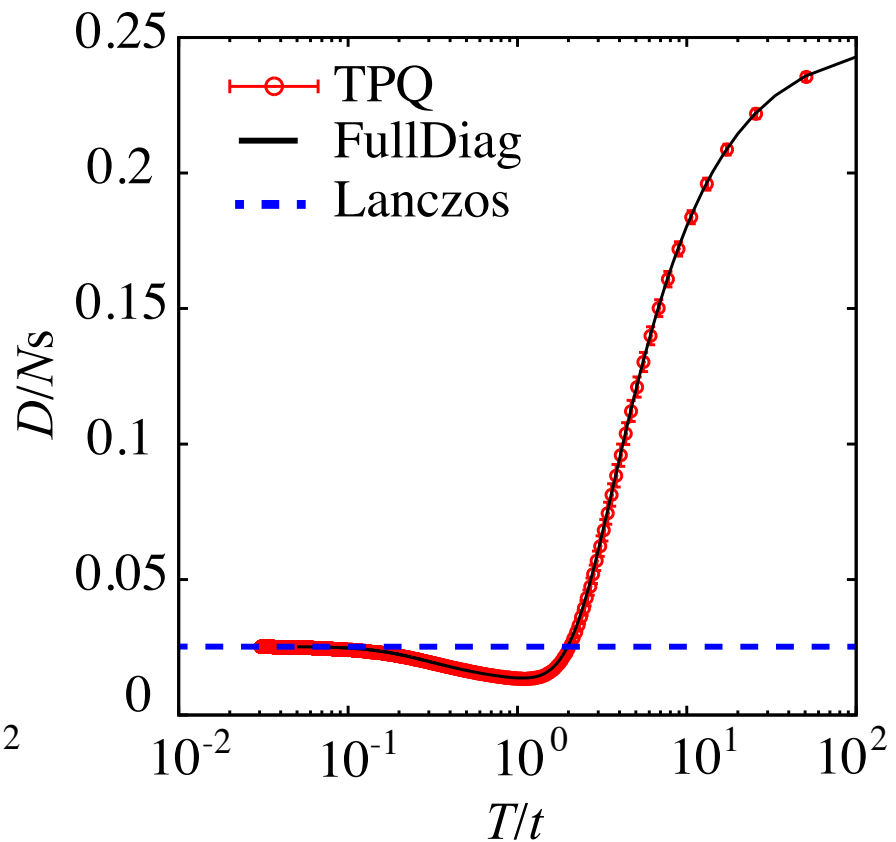
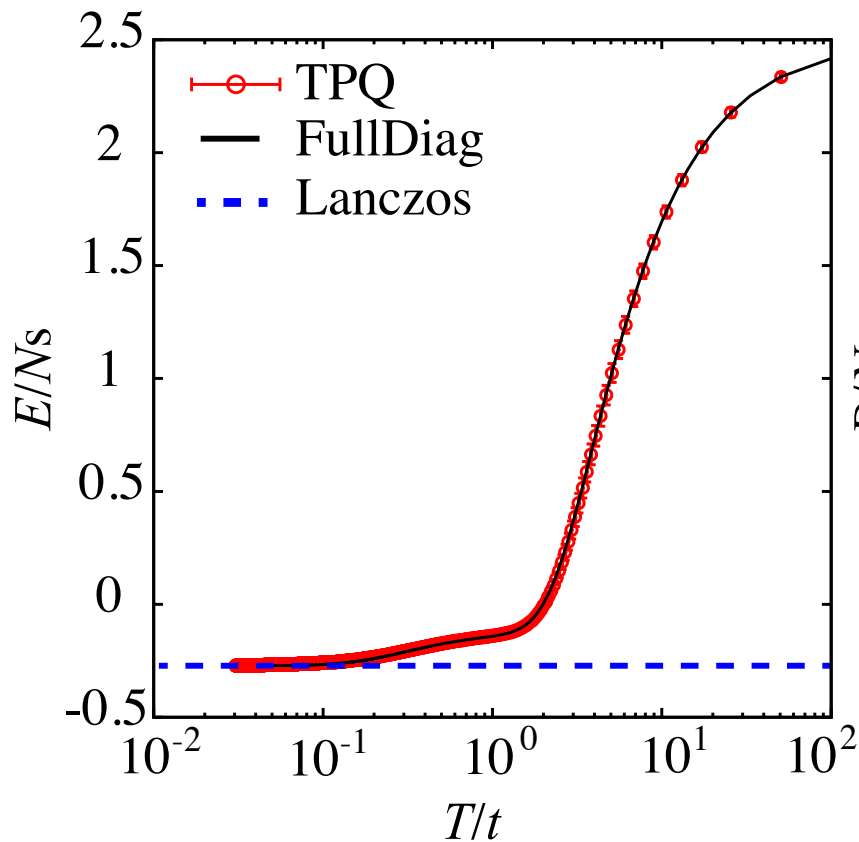
dimension of vector $\sim 10^8 \times 10^8$

Memory $\sim 10\text{ G Byte}$

\rightarrow Possible even in lab's cluster machine !

Validity of TPQ method

Comparison of FullDiag, TPQ, Lanczos method
Hubbard model, $L=8$, $U/t=8$, half filling, $S_z=0$



TPQ method works well !

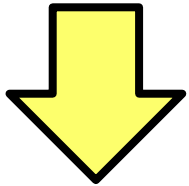


Let's start $H\Phi$!
(Standard mode)

How to use H Φ : Standard mode I (Lanczos)

Only **StdFace.def** is necessary (< 10 lines) !

```
L      = 4
model  = "Spin"
method = "Lanczos"
lattice = "square lattice"
J      = 1.0
2Sz   = 0
```



HPhi -s StdFace.def

./output : results are output

Important files

./output/zvo_energy.dat	→ energy
./output/zvo_Lanczos_Step.dat	→ convergence
./output/zvo_cisajs.dat	→ one-body Green func.
./output/zvo_cisajscktalt.dat	→ two-body Green func.

ex. 4×4 2d Heisenberg model,
GS by Lanczos method

Method

Lanczos — ground state

CG - LOBCG

TPQ — finite-temperature

FullDiag — full-diagonalization

How to use HΦ: Standard mode II

`./output/zvo_energy.dat`

```
$ cat output/zvo_energy.dat  
Energy -11.2284832084288109  
Doublon 0.000000000000000000  
Sz 0.000000000000000000
```

ex. 4by4, 2d Heisenberg model,
GS calculations by Lanczos

GS energy

`./output/zvo_Lanczos_Step.dat`

```
$ tail output/zvo_Lanczos_Step.dat  
stp=28 -11.2284832084 -9.5176841765 -8.7981539671 -8.5328120558  
stp=30 -11.2284832084 -9.5176875029 -8.8254961060 -8.7872255591  
stp=32 -11.2284832084 -9.5176879460 -8.8776934418 -8.7939798590  
stp=34 -11.2284832084 -9.5176879812 -8.8852955092 -8.7943260103  
stp=36 -11.2284832084 -9.5176879838 -8.8863380562 -8.7943736678  
stp=38 -11.2284832084 -9.5176879839 -8.8864307327 -8.7943782609  
stp=40 -11.2284832084 -9.5176879839 -8.8864405361 -8.7943787937  
stp=42 -11.2284832084 -9.5176879839 -8.8864422628 -8.7943788984  
stp=44 -11.2284832084 -9.5176879839 -8.8864424018 -8.7943789077  
stp=46 -11.2284832084 -9.5176879839 -8.8864424075 -8.7943789081
```

convergence process by Lanczos method

How to use HΦ: Standard mode III

`./output/zvo_cisajs.dat`

$$\langle c_{i\sigma}^\dagger c_{j\tau} \rangle$$

```
$ head output/zvo_cisajs.dat
```

```
0 0 0 0 0.50000000000 0.00000000000  $\langle c_{0\downarrow}^\dagger c_{0\downarrow} \rangle$ 
```

```
0 1 0 1 0.50000000000 0.00000000000  $\langle c_{0\uparrow}^\dagger c_{0\uparrow} \rangle$ 
```

`./output/zvo_cisajscktalt.dat`

```
$ head output/zvo_cisajscktalt.dat
```

```
0 0 0 0 0 0 0 0 0.50000000000 0.00000000000
0 0 0 0 0 1 0 1 0.00000000000 0.00000000000
0 0 0 0 1 0 1 0 0.1330366332 0.00000000000
0 0 0 0 1 1 1 1 0.3669633668 0.00000000000
```

$$\langle c_{0\downarrow}^\dagger c_{0\downarrow} c_{0\downarrow}^\dagger c_{0\downarrow} \rangle$$

$$\langle c_{0\downarrow}^\dagger c_{0\downarrow} c_{0\uparrow}^\dagger c_{0\uparrow} \rangle$$

$$\langle c_{0\downarrow}^\dagger c_{0\downarrow} c_{1\downarrow}^\dagger c_{1\downarrow} \rangle$$

$$\langle c_{0\downarrow}^\dagger c_{0\downarrow} c_{1\uparrow}^\dagger c_{1\uparrow} \rangle$$

ex. onsite • nn-site correlation func.

How to use HΦ: Standard mode IV

HPhi/samples/Standard/

StdFace.def for

Hubbard model, Heisenberg model, Kitaev model,
Kondo-lattice model

By changing StdFace.def slightly, you can easily perform the calculations for different models.

Cautions :

- Do not input **too large system size**

(upper limit@laptop: spin 1/2→24 sites, Hubbard model 12 sites)

- Lanczos method is unstable for **too small size**
(dim. > 1000)

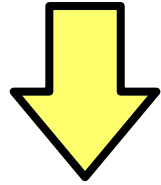
-TPQ method does not work well **for small size**
(dim. > 1000)



Expert mode !

How to use HΦ: What is Expert mode ?

HPhi -s StdFace.def



Standard mode: Necessary input files are **automatically generated**

Files for Hamiltonian

CoulombIntra.def, Trans.def, locspn.def...

Files for basic parameters (**two** files)

modpara.def, calcmod.def

Files for correlations functions (**two** files)

greenone.def, greentwo.def

+ list of input files: namelist.def

Expert mode: preparing the following files by yourself

How to use HΦ: What is Expert mode ?

Expert mode: preparing the following files by yourself

Files for Hamiltonian

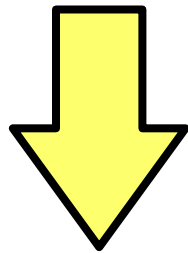
CoulombIntra.def, Trans.def, locspn.def...

Files for basic parameters (**two** files)

modpara.def, calcmod.def

Files for correlations functions (**two** files)

greenone.def, greentwo.def



execute following command

```
HPhi -e namelist.def
```



Enjoy HΦ



Questions ?