

今後の将来展望: HΦ

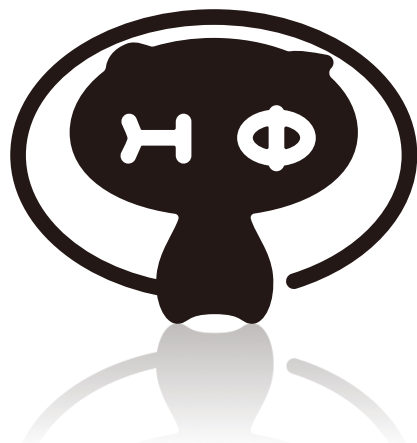
山地 洋平

東京大学大学院工学系物理工学専攻


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1. $K\omega$ の応用: Finite- T linear response
2. 今後の開発目標 Perspective



Computational
Science Alliance
The University of Tokyo



New function will be implemented: Finite- T linear response Combination of TPQ and $K\omega$

Y. Yamaji, T. Suzuki, & M. Kawamura, arXiv:1802.02854.



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Finite-Temperature Spectra

$$\mathcal{G}_\beta^{AB}(\omega) = \sum_{n,m} \frac{e^{-\beta E_n}}{Z(\beta)} \frac{\langle n | \hat{A}^\dagger | m \rangle \langle m | \hat{B} | n \rangle}{\omega + i\delta + E_n - E_m}$$

$$Z(\beta) = \sum_n e^{-\beta E_n}$$

$$\mathcal{G}_\beta^{AB}(\omega) = \sum_n \frac{e^{-\beta E_n}}{Z(\beta)} \langle n | \hat{A}^\dagger \frac{1}{\omega + i\delta + E_n - \hat{H}} \hat{B} | n \rangle$$

Complexity $\mathcal{O}(N_{\text{H}}^3)$

Memory $\mathcal{O}(N_{\text{H}}^2)$

Is it necessary? Answer is No

Finite-Temperature Spectra by Real-Time Evolution of Wave Functions

T. Iitaka and T. Ebisuzaki, Phys. Rev. Lett. 90, 047203 (2003).

R. Steinigeweg, J. Gemmer, and W. Brenig, Phys. Rev. Lett. 112, 120601 (2014).

T. Monnai and A. Sugita, J. Phys. Soc. Jpn. 83, 094001 (2014).

C. Karrasch, D. M. Kennes, and J. E. Moore, Phys. Rev. B 90, 155104 (2014).

F. Jin, R. Steinigeweg, F. Heidrich-Meisner, K. Michielsen, and H. De Raedt,
Phys. Rev. B 92, 205103 (2015).

Finite-Temperature Spectra by Microcanonical Ensemble

M. W. Long, P. Prelovsek, S. El Shawish, J. Karadamoglou, and X. Zotos,
Phys. Rev. B 68, 235106 (2003).

X. Zotos, Phys. Rev. Lett. 92, 067202 (2004).

An Intuitive Description of TPQ States and Green's Function at Finite Temperature

A normalized TPQ state

$$|\psi_\beta\rangle \equiv \frac{|\phi_\beta\rangle}{\sqrt{\langle\phi_\beta|\phi_\beta\rangle}} \sim \sum_n e^{i\varphi_n} \frac{e^{-\frac{\beta}{2}E_n}}{\sqrt{Z(\beta)}} |n\rangle$$

Spectral projector $\hat{P}_n = |n\rangle\langle n|$

Green's function rewritten by using a TPQ state

$$\mathcal{G}_\beta^{AB}(\zeta) \sim \sum_n \langle\psi_\beta|\hat{P}_n\hat{A}^\dagger \frac{1}{\zeta + E_n - \hat{H}} \hat{B}\hat{P}_n|\psi_\beta\rangle$$

An Alternative to Spectral Projection

T. Kato, Progress of Theoretical Physics 4, 514 (1949).

$$\hat{P}_{\gamma,\rho} = \frac{1}{2\pi i} \oint_{C_{\gamma,\rho}} \frac{dz}{z - \hat{H}} \quad z = \rho e^{i\theta} + \gamma$$

$$|\phi\rangle = \sum_n d_n |n\rangle$$

$$\hat{P}_{\gamma,\rho} |\phi\rangle = \sum_{E_n \in (\gamma - \rho, \gamma + \rho)} d_n |n\rangle$$

Discretized by Riemann sum

$$\hat{P}_{\gamma,\rho,M} = \frac{1}{M} \sum_{j=1}^M \frac{\rho e^{i\theta_j}}{\rho e^{i\theta_j} + \gamma - \hat{H}}$$

$$\theta_j = 2\pi(j - 1/2)/M$$

T. Sakurai and H. Sugiura,
J. Comput. Appl. Math. 159, 119 (2003).
T. Ikegami, T. Sakurai, and U. Nagashima,
J. Comput. Appl. Math. 233, 1927 (2010).

Shifted Krylov Subspace Method

$$\vec{x} = \frac{1}{\rho e^{i\theta_j} + \gamma - \hat{H}} \vec{b}$$

Liner equations

$$(z\mathbf{1} - H)\vec{x} = \vec{b}$$

$$\vec{b} \doteq \hat{O}|\psi\rangle$$

$$\Rightarrow G_{\hat{O}}(z) = \vec{b}^\dagger \vec{x}$$

$$\vec{x} \doteq (z\mathbf{1} - \hat{H})^{-1} \hat{O}|\psi\rangle$$

← Solvable by Shifted Krylov subspace method

A. Frommer (1995, 2003)

T. Sogabe, T. Hoshi, S. L. Zhang, and T. Fujiwara, *A numerical method for calculating the Green's function arising from electronic structure theory*, In *Frontiers of Computational Science*. pp.189-195, 2007.

Shifted CG: Algorithm

Initial $\vec{r}_0 = \vec{b}$, $\alpha_{-1} = 1$, $\rho_{-1} = +\infty$,
 $\pi_0^\sigma = \pi_{-1}^\sigma = 1$, $\vec{p}_{-1}^\sigma = \vec{x}_{-1}^\sigma = \vec{0}$

For $k = 0, 1, \dots, m$

-Seed equations

$$\rho_k = \vec{r}_k^T \vec{r}_k$$

$$\beta_{k-1} = \frac{\rho_k}{\rho_{k-1}}$$

$$\alpha_k = \frac{\rho_k}{\vec{r}_k^T A \vec{r}_k - \beta_{k-1} \frac{\rho_k}{\alpha_{k-1}}}$$

$$\vec{r}_{k+1} = \left(1 + \frac{\alpha_k \beta_{k-1}}{\alpha_{k-1}}\right) \vec{r}_k - \alpha_k A \vec{r}_k - \frac{\alpha_k \beta_{k-1}}{\alpha_{k-1}} \vec{r}_{k-1}$$

-Shifted equations

$$\pi_{k+1}^\sigma = (1 + \alpha_k \sigma) \pi_k^\sigma - \frac{\alpha_k \beta_{k-1}}{\alpha_{k-1}} (\pi_{k-1}^\sigma - \pi_k^\sigma)$$

$$\vec{p}_k^\sigma = \frac{1}{\pi_k^\sigma} \vec{r}_k + \beta_{k-1} \left(\frac{\pi_{k-1}^\sigma}{\pi_k^\sigma}\right)^2 \vec{p}_{k-1}^\sigma$$

$$\vec{x}_k^\sigma = \vec{x}_{k-1}^\sigma + \frac{\pi_k^\sigma}{\pi_{k+1}^\sigma} \alpha_k \vec{p}_k^\sigma$$

Seed switch

S. Yamamoto, *et al.*,
 J. Phys. Soc. Jpn. 77, 114713 (2008).

Library $K\omega$

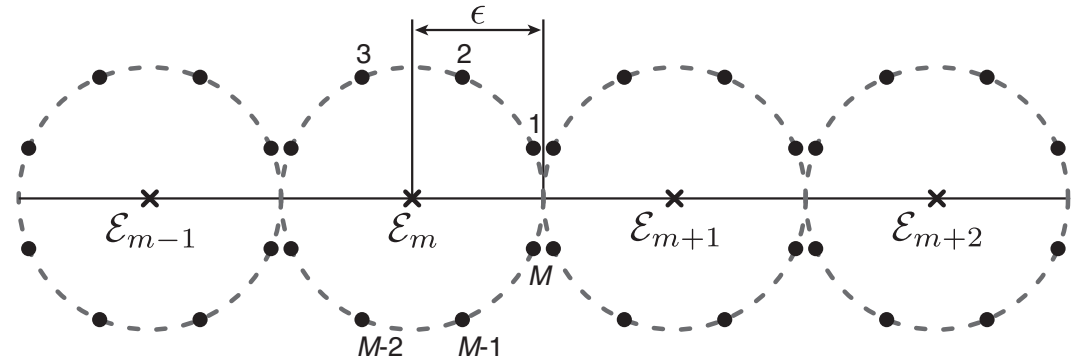
by Dr. Kawamura (ISSP)



Finite-Temperature Green's Function by Typical Pure States

$$|\psi_{\beta, \delta}^m\rangle = \hat{P}_{\mathcal{E}_m, \epsilon, M} |\psi_{\beta}\rangle$$

$$\delta = (E_0, \epsilon, M)$$



$$\mathcal{E}_m = E_0 + (2m - 1)\epsilon$$

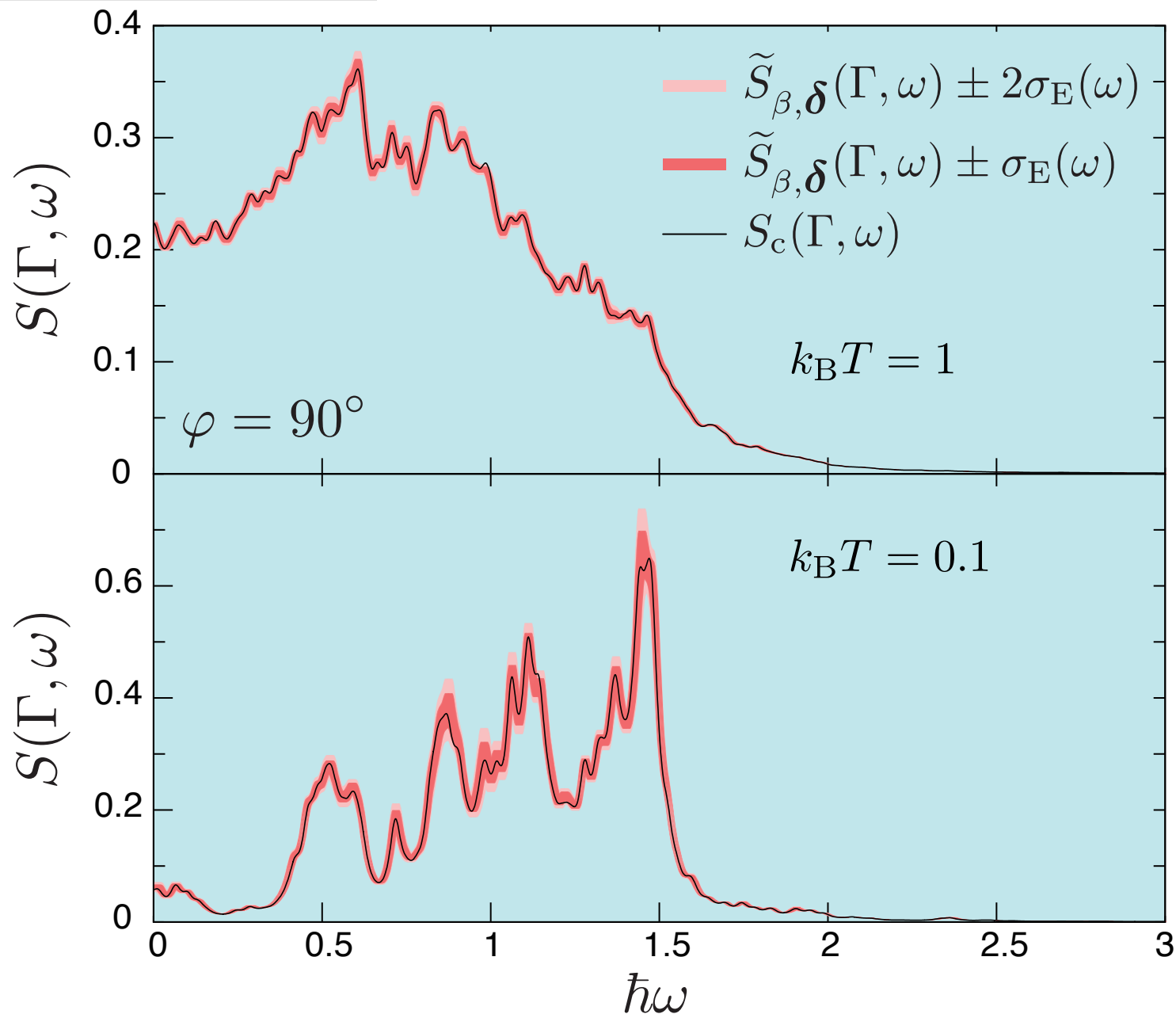
Green's function

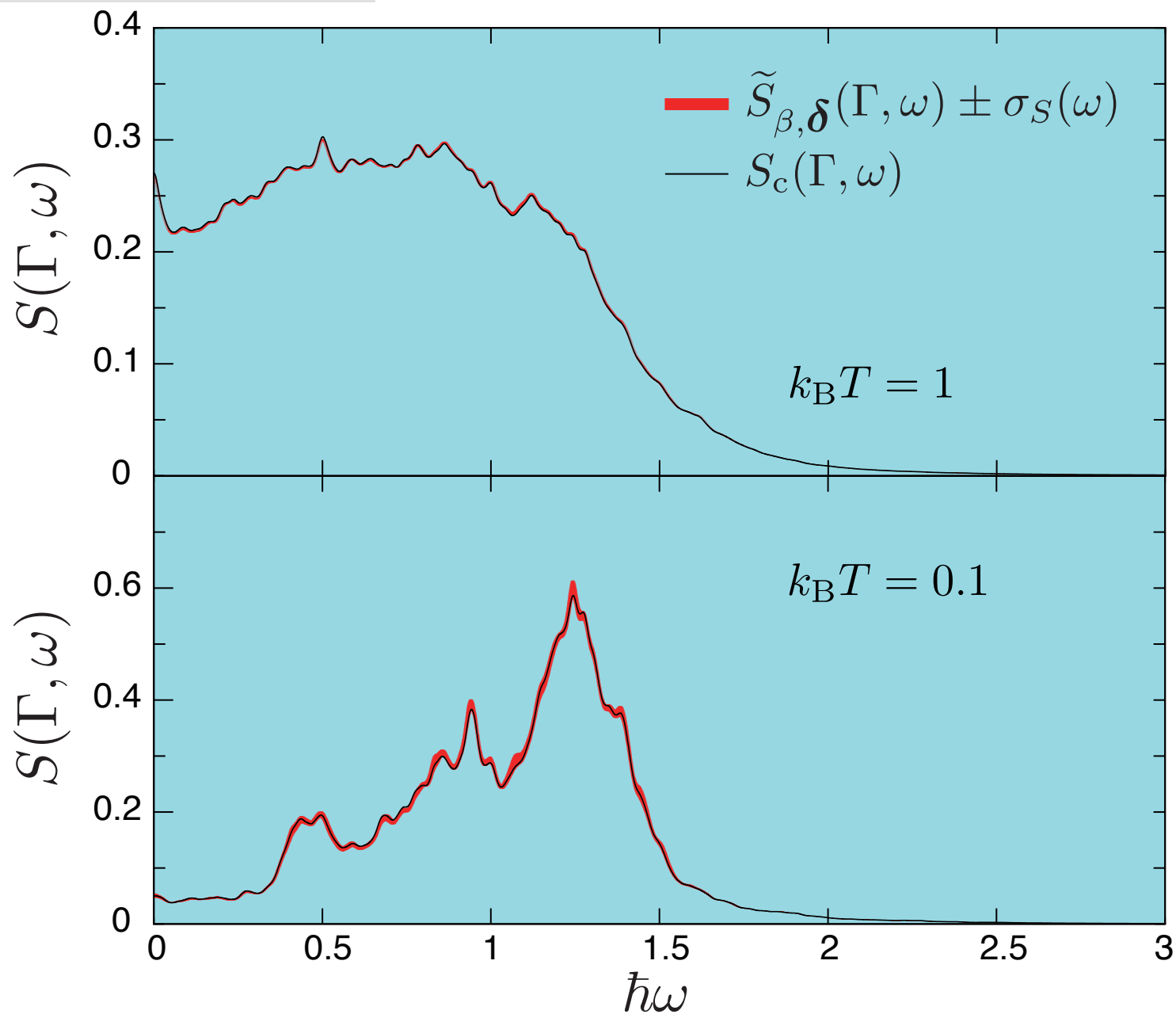
$$\tilde{\mathcal{G}}_{\beta, \delta}^{AB}(\zeta) = \sum_{m \geq 0} \langle \psi_{\beta, \delta}^m | \hat{A}^\dagger \frac{1}{\zeta + \mathcal{E}_m - \hat{H}} \hat{B} | \psi_{\beta, \delta}^m \rangle$$

$$\mathcal{G}_{\beta}^{AB}(\zeta) = \lim_{\epsilon \rightarrow +0} \lim_{M \rightarrow +\infty} \mathbb{E} \left[\tilde{\mathcal{G}}_{\beta, \delta}^{AB}(\zeta) \right]$$

Probability distribution

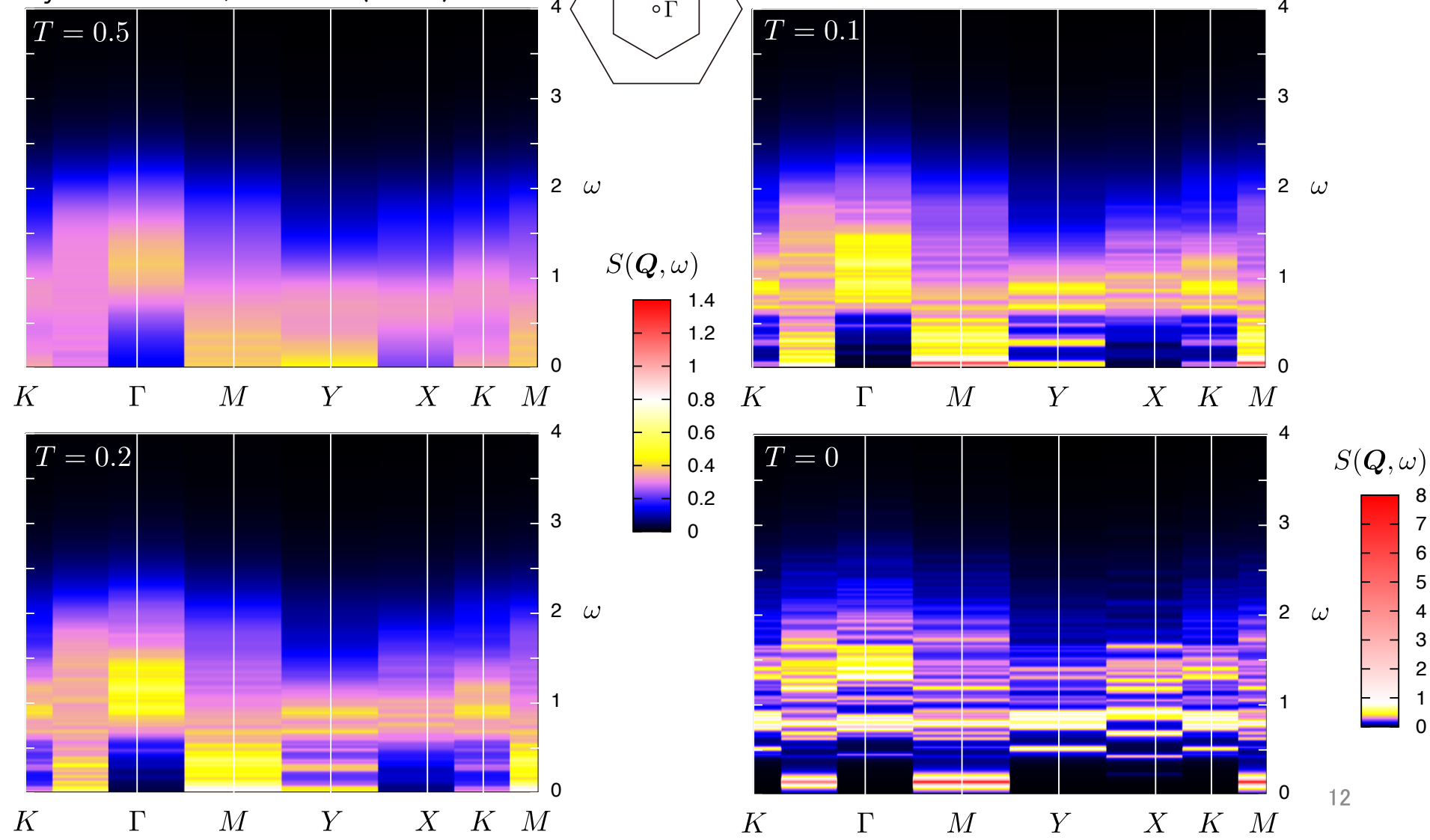
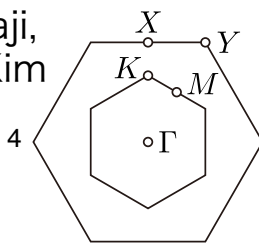
$$\tilde{P}_{\delta}(\mathcal{E}_m) = \langle \psi_{\beta, \delta}^m | \psi_{\beta, \delta}^m \rangle$$





Finite- T $S(Q, \omega)$ of a Frustrated Magnets: Γ model

A. M. Samarakoon, G. Wachtel, Y. Yamaji,
D. A. Tennant, C. D. Batista, and Y. B. Kim
Phys. Rev. B 98, 045121 (2018)



Future Plan

New functions will be implemented

1. Finite- T linear response:
Combination of TPQ and $K\omega$
2. N spin/body interaction and Green's function
3. Tool for optimizing model parameters to fit experimental measurements
-Example: Find an effective spin Hamiltonian that reproduces an observed magnetizaion process
4. Symmetry
-Reduction of dimension of Hilbert space
-Analysis of wave fucntions